Subdivision Surfaces: A New Paradigm For Thin-Shell Finite-Element Analysis

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- 1. The problem with thin-shell finite-element analysis
- 2. Physics of thin-shells
- 3. Finite element discretization
- 4. Subdivision surfaces
- 5. Examples, convergence of method
- 6. Conclusions
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# 1. Thin Shell Deformations The Problem

• Undeformed shell, deformed shell, finite element analysis



## **Thin-Shell Surface Problem**

- Difficult to create C<sup>1</sup> continuity between Elements of the limit surface of a shell
- Shell must have a finite Kirchhoff-Love energy
- Usual methods include derivatives, lead to high order polynomials, difficult to calculate and physical limitations
- Purpose of paper: present a method that leads to the desired C<sup>1</sup> continuity.

## 2. Physics of deformation in Thin-Shell

- Energies involved in a deformation:
  - Strain of the shell (elasticity)



#### - Variation of the bending of the shell



$$\alpha_{ij} = \frac{1}{2}(a_i a_j - a_i a_j)$$

Strain tensor

$$\beta_{\alpha\beta} = a_{\alpha} \frac{\partial a_{3}}{\partial \theta^{\beta}} - \overline{a}_{\alpha} \frac{\partial a_{3}}{\partial \theta^{\beta}}$$

Bending strain

where *a*, *a* are basis vectors in undeformed and deformed shell respectively  $i, j \in \{1,2,3\}, \alpha, \beta \in \{1,2\}$ are the 3D components  $(\theta^1, \theta^2, \theta^3)$ 

## Physics of deformation in Thin-Shell

Since the deformed shell is the undeformed shell plus a deformation function (*linearization*), i.e.

 $x(\theta^1, \theta^2) = \overline{x}(\theta^1, \theta^2) + u(\theta^1, \theta^2)$ 

where *x*, *x* are the coordinates in the deformed and the undeformed configuration, *u* is displacement function

... the two deformation tensors can be expressed as a function of the not deformed coordinates and the displacement functions, which will be ideal for the finite-element analysis.

## Equilibrium in the Deformation

- The energy density of the shell is a function of the previously defined α and β;
- The potential energy of the shell thus is

$$\Phi^{\rm int}[u] = \int_{\Omega} W(\alpha, \beta) d\Omega$$

• The potential energy of the applied load is  $\Phi^{ext}[u] = -\int_{\Omega} q \cdot u d\Omega - \int_{\partial\Omega} N \cdot u ds$ 

where *q* are applied loads, *N* axial forces on boundary

## Equilibrium in the Deformation

- In a stable configuration the sum of the potential energies must be minimal (physics)
- Potential energy:  $\Phi[u] = \Phi^{int}[u] + \Phi^{ext}[u]$ We minimize it according to Euler-Lagrange equations

 $\langle D\Phi[u], \delta u \rangle = \langle D\Phi^{int}[u], \delta u \rangle + \langle D\Phi^{ext}[u], \delta u \rangle = 0$ 

 "Statement of the principle of virtual work": Actio=Reactio principle, force caused by shell deformation must be compensated by the force caused by the loads.

## **3.** Finite-Element Discretization

- Build a mesh on the shell, choose base function
- The discretization leads to the expression:

K is the energy of the shell

u is displacement field (array)

f is external force applied to the shell

• Done as sum over the elements

 $K_h \cdot u_h = f_h$ 

- K and f involve the evaluation of an integral
- Integrals can be computed with a quadrature rule
- Authors use a one-point quadrature rule, which is said to achive a sufficient precision for this analysis.

## 4. Subdivision Surfaces

- Construction of a smooth surface
- Done by repeated subdivision of a given mesh
- New nodes created at every subdivision
- Coordinates of nodes at step k+1 are computed as linear combination of nodes at step k
- Good choice of weights produce a smooth limit surface

(H<sup>2</sup> integrability, C<sup>1</sup> continuity)

## Subdivision Schemes: Why?

- 1D case easy to build C<sup>n</sup> continuity through polynomial interpolation
- 2D surfaces: C<sup>2</sup> smoothness requires up to 6th order polynomials
- Difficulties arise when handling cross-patch smoothness
- Approximation scheme: C<sup>2</sup> continuity
  => subdivision surfaces are advantageous!

## Subdivision Scheme: Loop

- Subdivision for triangulated meshes done with Loop's scheme, although every strategy could be used
- Leads to quadrisection of every triangle



### Loop's Scheme



$$x_{0}^{k+1} = (1 - Nw)x_{0}^{k} + wx_{1}^{k} + \dots + wx_{N}^{k}$$
  
where  
$$w = \frac{1}{N} \left[ \frac{5}{8} - \left( \frac{3}{8} + \frac{1}{4} \cos(\frac{2\pi}{N})^{2} \right] \text{(Loop)} \right]$$
$$w = \begin{cases} \frac{3}{8N}, N > 3\\ \frac{3}{16}, N = 3 \end{cases} \text{(Warren's Choice)} \end{cases}$$



$$x_{I}^{k+1} = \frac{3x_{0}^{k} + x_{I-1}^{k} + 3x_{I}^{k} + x_{I+1}^{k}}{8}$$

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#### Loop's Scheme: Examples



#### Loop on a Distributor Cap mesh

#### Loop's Scheme: Examples





### **Convergence of limit surface**

- The convergence of this method can be proven using a two vertex neighbourhood.
- Calculation of limit position of vertices using a one vertex neighbourhood:

Let  $x^k = (x_0^k, x_1^k, ..., x_N^k)$ , where N is valence of vertex. We can express  $x^{k+1}$  as  $x^{k+1} = S \cdot x^k$ 

S is a matrix expressing the Loop relationship. Computation of limit configuration of vertices  $(k \rightarrow \infty)$ 

$$x^{\infty} = (S)^{\infty} \cdot x^0$$
,  $x^0$  start configuration

#### **Convergence of limit surface**

 $x^{\infty} = (S)^{\infty} \cdot x^0$ 

Since  $\lambda_0=1$ ,  $\lambda_i\leq 1$  for all  $i\neq 0$ , S<sup> $\infty$ </sup> converges to a limit.

Using eigenvalue/eigenvector decomposition, it can be shown that :

$$X^{\infty} = L_0 \cdot X^0$$



where  $L_0 = (1 - N \cdot l, l, ..., l),$   $l = \frac{1}{\frac{3}{8w} + N}$ 

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## **Convergence of limit surface**

- Similar consideration for all the limit values needed in the FE computation (tangents on shell, surface normal)
- Advantage: computation of limit configuration of vertices and other primitives is possible at every step *k* of the refinement, i.e. when one has achieved the desired mesh subdivision.
- Convergence (regular patches, i.e. valence 6 at each vertex) to a quartic box Spline on every patch.

#### **Evaluations in an Element**

- We need to compute values of points and derivates inside an element (FE Analysis).
- For regular patches (vertices with valence 6):



$$x(\theta^{1}, \theta^{2}) = \sum_{i=1}^{12} N^{i}(\theta^{1}, \theta^{2}) \cdot x_{i}$$
$$u_{h}(\theta^{1}, \theta^{2}) = \sum_{i=1}^{12} N^{i}(\theta^{1}, \theta^{2}) \cdot u_{i}$$
$$M \text{ spling shape functions}$$

x, u vertex coordinates and displacements of neighbourhood

## **Evaluations in an Element**

 Irregular patches: subdivide the element with Loop's scheme until the searched point is known to be in a regular sub-patch, then compute as before (with adapted parameters)





#### Implementation and computation

- 1. One subdivision step (Max one irregular vertex per patch)
- 2. Introduction of artificial nodes at boundary
- 3. Find 1-neighbourhood of vertices
- 4. Create local coordinates on irregular patches
- 5. Create stiffness matrix and force array
- 6. Introduce displacement boundaries
- 7. Solve system of equations (finite elements)
- 8. Compute limit position of nodes (sub. surfaces)

## 5. Examples and convergence

- The method is compared with two other approaches.
- A bound for a finite-element solution is known to exist.
- For the examples shown in the following, an analytical solution is known, thus we can analyse the "goodness" of this approach with the exact solution as well.

## **Rectangular Plate**



A typical mesh on such a plate Irregular vertices are present Uniform load on the shell



Clamped Boundary Simply Supported Boundary

## **Pinched Cylinder**







- Due to subdivision scheme the loads spread over several points.
- The total weight is maintained

## **Pinched Cylinder**



- Method converges to the optimal solution
- Convergence is faster than two other methods

## **Hemispherical Shell**



- Surface cannot be triangulated without irregular nodes
- Generalizations needed



 Hard Test : Generic Box-Spline Approach not possible

#### **Hemispherical Shell**



- Surface converges optimally for this irregular mesh as well, even if standard approach is not possible
- Important that no parasitic strains appear.

## 6. Conclusions

- Use of subdivision surfaces for description of undeformed and deformed shell
- Method takes care of physical considerations (finite Kirchhoff-Love energy)
- Loop scheme: provable local convergence
- Smoothness between elements without using derivatives
- Finite element analysis on same mesh as subdivision (no additional triangulation error)

## Conclusions

- Displacement field depends not only from element vertices, but from the 1-Neighbourhood as well
- Simple quadrature for finite-elements is sufficient
- Convergence is optimal in the finite element sense
- Method is applicable as well for other subdivision rules, not only for Loop scheme

# 7. My Opinion

- Easy to implement (easier than considering derivatives, mask existence, no particular data structures)
- No double meshes needed
- Respects physical laws

## It would be interesting to see...

- How other schemes really behave
- Behaviour with not linearized Kinematics







#### Kirchhoff-Love Energy

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